

# **Facile hydrothermal synthesis and optical limiting properties of TiO<sub>2</sub>-reduced graphene oxide nanocomposites**

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**Abstract:**  $\text{TiO}_2$ /reduced graphene oxide (RGO) nanocomposites  $\text{G}_x$  (RGO Titania nanocomposite, x grams tetrabutyl titanate per 0.03 grams RGO,  $x = 0.25, 0.50, 1.00$ ) were prepared by a hydrothermal method: graphene oxide was reduced to RGO in a 2:1 water:ethanol mixture in the presence of varying quantities of tetrabutyl titanate, which deposited as  $\text{TiO}_2$  on the RGO sheets. The nanocomposites were characterized by a combination of Fourier transform infrared spectroscopy, diffuse reflectance ultraviolet-visible spectroscopy, photoluminescence spectroscopy, Raman spectroscopy, X-ray powder diffraction, X-ray photoelectron spectroscopy and transmission electron microscopy studies. The nanocomposite  $\text{G}_{0.25}$  exhibits enhanced nonlinear optical properties compared to its individual components, which is ascribed to a combination of mechanisms. The role of defects and electron/energy transfer in the optical limiting performance of  $\text{G}_{0.25}$  was clarified with the help of Raman and photoluminescence spectroscopies. Intensity-dependent switching between reverse saturable absorption and saturable absorption behavior was observed with the  $\text{G}_{0.50}$  nanocomposite.

## 1. Introduction

Graphene, a two-dimensional  $sp^2$ -hybridized carbon nanosheet, possesses an atomically thin planar geometry, an unusual band structure and a high carrier mobility, and has spurred immense interest in novel graphene-based materials for a variety of technological applications such as nanocomposites, energy-related materials and optoelectronic devices [1-4]. The majority of work on graphene nanocomposites has been based on reduced graphene oxide (RGO) because of the excellent scalability of its synthesis and fact that defect-free graphene sheets have recently become available. Compared to graphene sheets, RGO nanosheets are pock-marked with defects [5]. In contrast to graphene oxide (GO, an electrically insulating material because of the disrupted  $sp^2$  bonding networks), RGO possesses electrical conductivity and electron/hole transporting properties as the  $\pi$  network is restored, which is favorable for the construction of optoelectronic devices [6]. In particular, significant absorption of incident light and ultrafast carrier dynamics per layer result in RGO being a fast absorber over a wide spectral range [7].

Titanium dioxide ( $TiO_2$ ), a semiconductor material, has been at the forefront of applied research in the fields of optically active coatings, photo-catalysis and energy conversion due to its specific optical and electronic properties, chemical stability, high optical transparency and large refractive index [8,9]. In comparison with pristine  $TiO_2$ , nanoscale composite materials containing  $TiO_2$  are more interesting because of their potential applications in optoelectronic and photovoltaic devices [10-12]. Some investigations have indicated that a doping procedure can inhibit the recombination of electron-hole pairs, as this introduces new energy levels in the bandgap, and thus accelerates the transfer of photogenerated electrons [13]. Although some efforts have been made to synthesize  $TiO_2$ -based nanocomposites and to investigate their linear optical properties [14,15], the nonlinear optical (NLO) performance of  $TiO_2$ -based nanocomposites remains largely unexploited thus far [16].

The advantages of RGO and the attractive merits of  $TiO_2$  suggest that the combination of the two classes of material may lead to new nanocomposite materials that possess novel properties. Indeed, previous studies of RGO/ $TiO_2$  materials have

highlighted the interesting properties of nanocomposites obtained by a variety of approaches [17-19]. In this study, we report the facile hydrothermal syntheses of TiO<sub>2</sub>-decorated RGO nanocomposites with varying composition, characterization of the resultant assemblies by a combination of Fourier transform infrared spectroscopy, diffuse reflectance ultraviolet-visible spectroscopy, photoluminescence spectroscopy, Raman spectroscopy, X-ray powder diffraction, X-ray photoelectron spectroscopy and transmission electron microscopy studies, and a systematic evaluation of the influence of the TiO<sub>2</sub> content on the NLO and optical limiting (OL) behavior of the TiO<sub>2</sub>/RGO nanocomposites.

## **2. Experimental**

### **2.1 Materials**

Purified natural graphite was purchased from Qingdao Zhongtian Co. Ltd. Tetrabutyl titanate and ethanol were supplied by Shanghai Sinopharm Chemical Reagent Co. Ltd. All other reagents were purchased from commercial suppliers and were used without further purification, unless otherwise stated. Distilled water was used for all experiments.

### **2.2 Instruments and measurements**

Fourier-transform infrared (FTIR) spectra were carried out with a MB154S-FTIR spectrometer (Bomem, Canada) using spectroscopic grade KBr pellets in the range 4000-400 cm<sup>-1</sup>; spectra were recorded by accumulating 64 scans at a 4 cm<sup>-1</sup> spectral resolution at room temperature. Raman spectra were recorded at room temperature on a Renishaw inVia Raman Microscope, exciting at 532 nm with an Ar<sup>+</sup> laser. The laser light was focused onto samples by using a microscope equipped with an a×100 objective. UV-Vis spectra were recorded using a Varian Cary 500 spectrophotometer equipped with a diffuse reflectance accessory at room temperature in the range 200-800 nm. A BaSO<sub>4</sub> pellet was used as a reflectance standard. Fluorescence spectra were measured using a Fluoro-Max-P spectrofluorimeter. Transmission electron microscopy (TEM) experiments were performed with a JEM-2100 (JEOL), working at 200 kV. Samples for TEM analysis were obtained by spreading a drop of a dilute dispersion of the as-prepared products on amorphous carbon-coated copper grids,

which were then dried in air before transfer to the TEM sample chamber. X-ray powder diffraction (XRD) experiments were carried out at room temperature on a XD-3 diffractometer (Beijing Purkinje General Instrument Co. Ltd, China) by using Cu K $\alpha$  radiation ( $\lambda = 0.15418$  nm). X-ray photoelectron spectroscopy (XPS) characterization was performed on a RBD upgraded PHI-5000C ESCA (PerkinElmer) electron spectrometer with a Mg K $\alpha$  line at 280 eV.

The Z-scan and OL measurement setup used in our study has been described before [20,21]. Linearly polarized 4 ns pulsed laser radiation at 532 nm was used for excitation. For measurements of the NLO properties, all samples were placed in 2 mm quartz cells, and adjusted to have the same transmittance. To avoid cumulative thermal effects interfering with experiments directed at identifying the major OL mechanism, the laser pulses were delivered at a low frequency of 2 Hz; each pulse of light was therefore ensured to encounter fresh sample.

### 2.3 Preparation of TiO<sub>2</sub>-decorated RGO nanocomposites

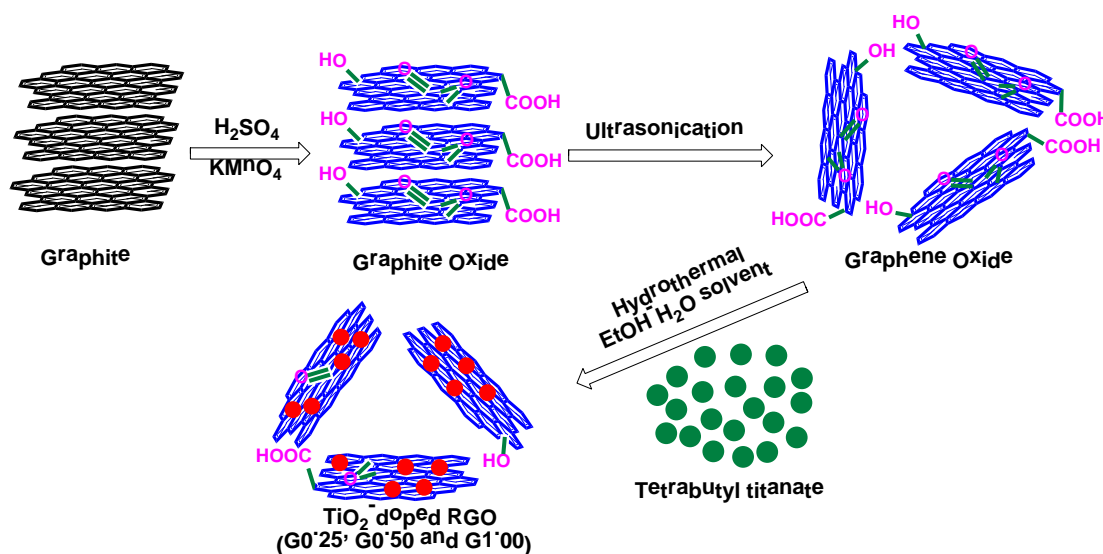


Figure 1. Illustration of the fabrication process for TiO<sub>2</sub>-decorated RGO nanocomposites.

GO prepared by a modified method of Hummers and Offeman [22] was used as a substrate for TiO<sub>2</sub> growth, affording TiO<sub>2</sub>/RGO nanocomposites. A simple hydrothermal method was used to prepare the nanocomposites, using EtOH-water as

solvent as outlined in Figure 1. Under such solvothermal conditions, EtOH-water can reduce GO to RGO [23,24]. In a typical synthesis, GO (50 mg) was first dissolved in a mixed solution of deionized H<sub>2</sub>O (6 mL) and EtOH (3 mL) by ultrasonic treatment for 30 min, resulting in a yellow-brown solution. Tetrabutyl titanate was then added to the GO solution, and the resultant mixture was stirred for a further 30 min to afford a homogeneous suspension. The suspension was placed in a 10 mL Teflon-sealed autoclave that was maintained at 160 °C for 72 h. During the hydrothermal process, GO was reduced to RGO with simultaneous deposition of TiO<sub>2</sub> onto the RGO sheet. The resultant composite was collected by centrifugation, washed with distilled water several times followed by a rinsing in EtOH, and then dried under vacuum at room temperature for 18 h. To investigate the effect of TiO<sub>2</sub> content on the NLO properties of the RGO/TiO<sub>2</sub> composites, the weight of tetrabutyl titanate added (as the precursor to TiO<sub>2</sub>) was varied (0.25, 0.50, or 1.00 g), the resulting samples being denoted G<sub>x</sub>, where  $x = 0.25, 0.50$  and  $1.00$ , respectively. For comparison, blank RGO and pure TiO<sub>2</sub> were prepared via a hydrothermal route under the same conditions, but without the addition of tetrabutyl titanate or GO, respectively.

### 3. Results and discussion

#### 3.1 Synthesis

Given the interest in the chemistry of RGO nanocomposites, it is desirable to have convenient dispersion methods to increase the processability and availability of RGO. In this study, a series of RGO/TiO<sub>2</sub> nanocomposites (denoted as G<sub>x</sub>, where  $x$  represents the amount of tetrabutyl titanate in grams, per 0.03 g of RGO) were prepared by a simple hydrothermal treatment of tetrabutyl titanate and GO in mixed EtOH-H<sub>2</sub>O. Under the solvothermal conditions, this results in the facile reduction of GO to RGO, while the nanoparticles of TiO<sub>2</sub> are simultaneously dispersed on the RGO sheets. The fabrication process for the TiO<sub>2</sub>-decorated RGO nanocomposites is shown in Figure 1. Stacked graphite is first oxidized to graphite oxide by KMnO<sub>4</sub> and 98% H<sub>2</sub>SO<sub>4</sub>, the graphite oxide containing some oxygen-containing functional groups, such as carboxyl, epoxide and hydroxide. The functional groups on the GO provide reactive anchoring sites for nucleation and growth of nanoparticles. A stable GO

suspension was then obtained by dispersing the graphite oxide in distilled water and EtOH with ultrasonication. The mixed solvent is crucial in controlling the rate of hydrolysis of tetrabutyl titanate. TiO<sub>2</sub>/RGO nanocomposites, with different weight addition ratios of tetrabutyl titanate, were then obtained via a hydrothermal method based on a modification of Zhang's procedure [24]. During this process, the hydrothermal temperature was fixed at 160 °C, facilitating crystallization of the coating materials on RGO in the anatase phase, and consistent with the approach of Liang et al. [25]. This synthetic strategy for preparing TiO<sub>2</sub>-RGO nanocomposites provides several advantages [26,27]: (1) the process is facile, efficient, safe and "green" since it does not use hazardous materials; (2) the closed system at relatively high temperature and internal pressure promotes the recovery of  $\pi$ -conjugation; (3) RGO nanosheets agglomerate and precipitate because of  $\pi$ - $\pi$  interactions between the sp<sup>2</sup> carbons of RGO sheets after reduction. The problems associated with aggregation of RGO nanosheets in solvents are addressed by introducing direct interactions between the sp<sup>2</sup> carbons of the RGO nanosheets and the TiO<sub>2</sub> precursors; (4) the TiO<sub>2</sub>-stabilized RGO solutions display good solution processability, which is favorable for the measurement of photophysical properties; (5) the electron/energy transport properties of RGO nanosheets are improved because of the charge trap screening and doping effects of TiO<sub>2</sub>.

### 3.2 FTIR spectra analysis

Confirmation of TiO<sub>2</sub> functionalization of the RGO sheets was obtained by FTIR spectroscopy. Figure 2 displays typical FTIR spectra of TiO<sub>2</sub>, GO, RGO and the nanocomposite materials G0.25, G0.50, and G1.00. In the FTIR spectrum of GO, the most characteristic features are the absorption bands at 3367, 1726, 1363, and 1069 cm<sup>-1</sup>, corresponding to the O-H groups linked to the basal plane of GO, the C=O stretching vibration of the carboxylic groups, the C-OH stretching vibration, and the C-O stretching vibration, respectively [28]. All of these bands, which are associated with the oxygen-containing functional groups, decrease dramatically in intensity or even disappear after solvothermal reduction, consistent with conversion of GO into RGO. However, a small fraction of the oxygen-containing functional groups persists

and facilitates the hybridizing of the TiO<sub>2</sub> nanoparticles on the RGO sheets. Pure TiO<sub>2</sub> displays a broad IR band between 400 and 1000 cm<sup>-1</sup>, corresponding to the vibration of the Ti-O-Ti bonds (Figure 2) [29]. An IR band is observed below 1000 cm<sup>-1</sup> for the TiO<sub>2</sub>/RGO nanocomposites; compared to the corresponding band in TiO<sub>2</sub>, it is both broader and shifted to higher wavenumber. This band is presumably a combination of Ti-O-Ti vibration and Ti-O-C vibration modes resulting from the chemical interaction of TiO<sub>2</sub> with RGO [30-32]. Deconvolution of the FTIR spectral traces for G0.25, G0.50 and G1.00 in the range from 1080 to 400 cm<sup>-1</sup> has been effected (inset of Figure 2), confirming the presence of these Ti-O stretching vibrations. The presence of Ti-O-C bands suggests that, during the hydrothermal treatment, the oxygen-containing functional groups of GO interact strongly with the surface hydroxyl groups of the TiO<sub>2</sub> nanoparticles to ensure that the nanoparticles remain attached, even after chemical cleaning and ultrasonication, to eventually afford the chemically-bonded TiO<sub>2</sub>/RGO nanocomposites [33]. An absorption band at 1630 cm<sup>-1</sup> is due to the skeletal vibration of the graphene sheets [23,24]; this peak is also observed in the FTIR spectrum of RGO prepared by hydrothermal reduction of GO. These IR results are consistent with the reduction of GO to RGO, and with grafting of the TiO<sub>2</sub> nanoparticles onto the RGO plane, but do not provide information about different physicochemical properties of the TiO<sub>2</sub>/RGO nanocomposites resulting from different addition ratios of TiO<sub>2</sub>.

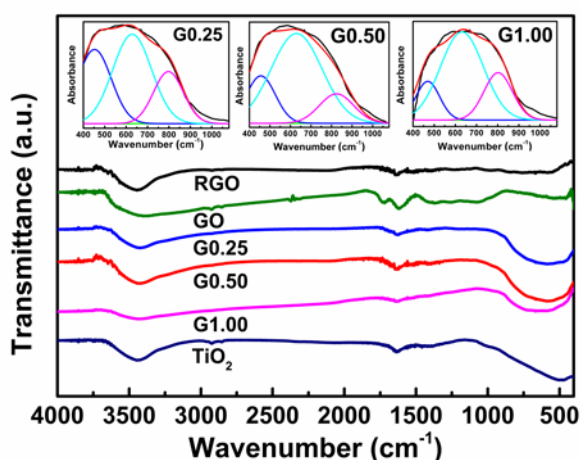
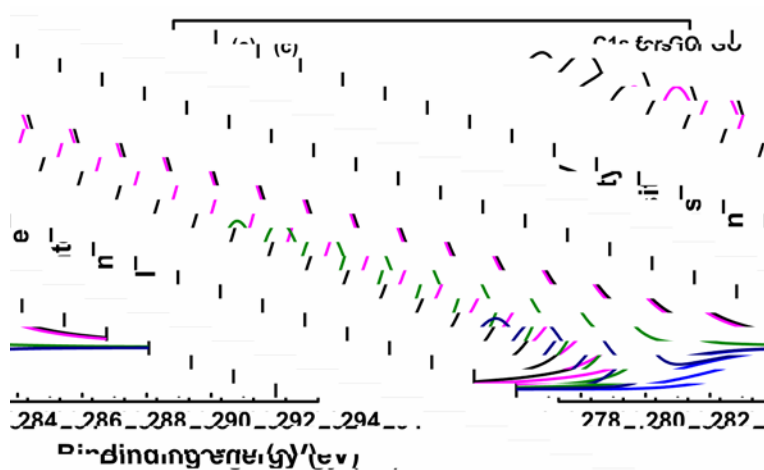
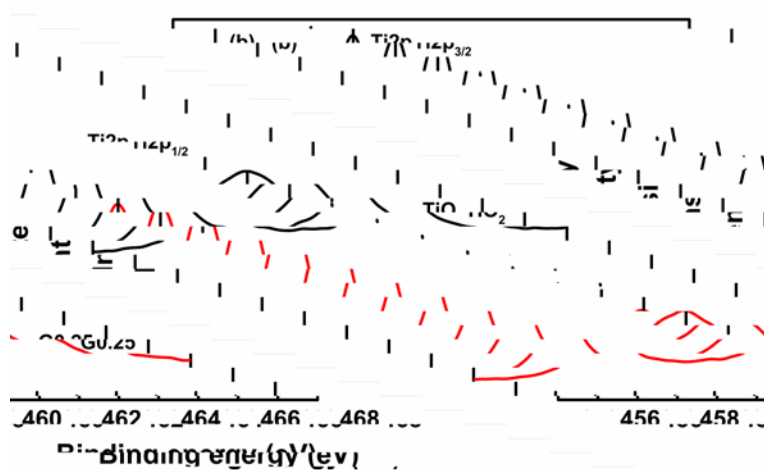
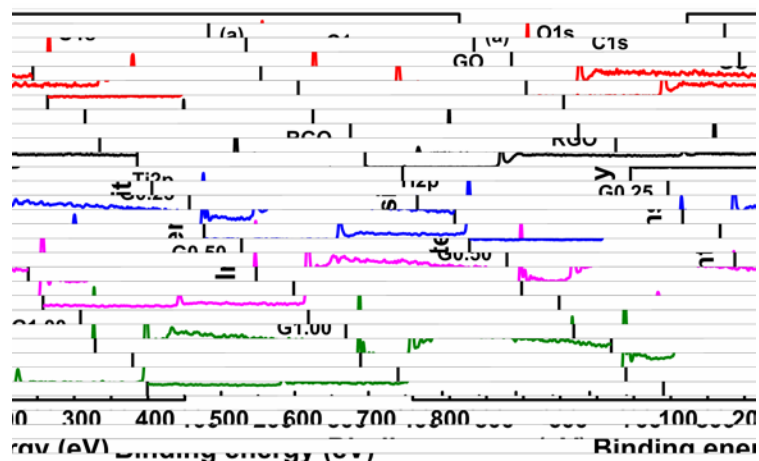


Figure 2. FTIR spectra of TiO<sub>2</sub>, GO, RGO, G0.25, G0.50 and G1.00. Shown in the



inset is the result of deconvoluting the spectra of G0.25, G0.50 and G1.00.

### 3.3 XPS spectra analysis



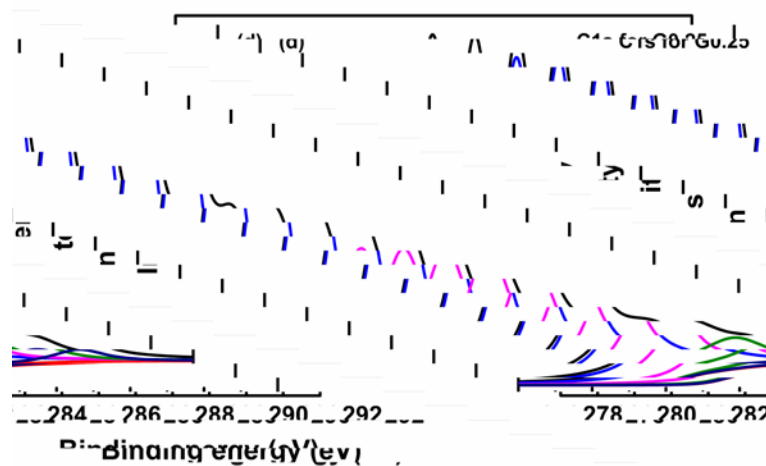


Figure 3. (a) XPS survey spectra of GO, RGO, G0.25, G0.50 and G1.00; (b) Ti 2p core level XPS spectra of TiO<sub>2</sub> and G0.25; High-resolution XPS spectra of C 1s for (c) GO and (d) the G0.25 nanocomposite.

The TiO<sub>2</sub>/RGO nanocomposites were characterized by XPS. From Figure 3(a), it can be seen that after thermal reduction, RGO shows a very weak O 1s signal which is much smaller than that of GO. This implies that a trace of residual oxygen functionalities still exists in the structure of RGO. For the TiO<sub>2</sub>/RGO nanocomposites, a new peak corresponding to Ti 2p can be clearly identified, with the binding energies at approximately 458.00, 459.00 and 457.00 eV for G0.25, G0.50 and G1.00, respectively. Figure 3(b) displays the Ti core XPS spectra of as-synthesized TiO<sub>2</sub> and the G0.25 nanocomposite. The Ti core-level XPS spectrum of TiO<sub>2</sub> exhibits two peaks centered at 464.25 and 458.50 eV, corresponding respectively to the Ti 2p<sub>1/2</sub> and Ti 2p<sub>3/2</sub> photoelectrons in the Ti<sup>4+</sup> state [34]. For the G0.25 nanocomposite, these two peaks occur at 464.58 and 458.83 eV, respectively. In each sample, the splitting between the two Ti-bands is about 5.75 eV, consistent with the presence of Ti<sup>4+</sup> in the TiO<sub>2</sub> and G0.25 nanocomposite [35]. The difference between the positions of the Ti peak of pure TiO<sub>2</sub> and that of the G0.25 nanocomposite may be assigned to the interactions of Ti with the oxygen centers of RGO. To investigate the carbon states in the nanocomposites, the high-resolution C 1s core levels of GO and G0.25 have been measured, as shown in Figures 3(c) and 3(d). Three typical peaks of GO observed at 283.90, 285.10 and 287.00 eV are usually ascribed to C=C and C-C bonds from GO,

and the oxygen-containing carbonaceous bands, respectively [6]. For the G0.25 nanocomposite, the asymmetrical and broad features of the observed C 1s XPS spectrum indicates the co-existence of distinguishable models. Deconvolution of the C 1s peaks shows the presence of chemically different C species in the TiO<sub>2</sub> matrix, with their binding energies at approximately 283.80, 285.20, 287.00, and 288.70 eV. The main C 1s peak located at 283.80 eV is attributed to defect-containing sp<sup>2</sup>-hybridized carbons. The sharp peak located at 285.20 eV is assigned to the presence of C-C bonds in RGO. In addition, there is a relatively weak peak located at 288.70 eV, contributed by carboxyl carbon (O=C-O) [36]. Such a surface functional group suggests that the -OH groups on the TiO<sub>2</sub> nanoparticles may react with the -COOH groups on the GO surface through esterification to form O=C-O-Ti bonds [37-39]. Moreover, the intensity of the peak at around 287.0 eV for the oxygen-containing groups decreases markedly and a notable enhancement of the sp<sup>2</sup>-C signal intensity corresponding to the C=C bond was observed in the C 1s XPS on proceeding to the G0.25 nanocomposite, implying that the introduced GO has been efficiently converted to RGO by thermal reduction. The peak ascribed to Ti-C bonds at around 282.00 eV is not detected in the G0.25 nanocomposite, indicating that carbon was not doped into the lattice of TiO<sub>2</sub> due to the low synthesis temperature. Similar results are also observed for G0.50 and G1.00 as are seen with G0.25, the results being shown in Figure S3.

### 3.4 UV-Vis spectra analysis

The UV-Vis diffuse reflectance spectra of the G0.25, G0.50 and G1.00 samples were carried out at room temperature; for comparison, the optical properties of TiO<sub>2</sub> nanoparticles were also measured under the same conditions. As shown in Figure 4(a), all samples exhibit some absorption in the visible region, with a  $\lambda_{\text{onset}}$  at 380-400 nm. The absorption of TiO<sub>2</sub> at ca. 400 nm is ascribed to charge transfer from the valence band (mainly formed from 2p orbitals of the oxide anions) to the conduction band (mainly formed from 3d t<sub>2g</sub> orbitals of the Ti<sup>4+</sup> cations) [40]. The presence of the varying amounts of TiO<sub>2</sub> significantly affects the light absorption of the TiO<sub>2</sub>/RGO nanocomposites; light absorption intensity in the UV region is significantly decreased

upon doping for all of the  $\text{TiO}_2/\text{RGO}$  nanocomposites with different addition ratios of  $\text{TiO}_2$ . Fan et al. [41] and Kim et al. [42] both observed a red-shift of the absorption edge of  $\text{TiO}_2$  and consequently a narrowing of the band gap of  $\text{TiO}_2$  upon doping with RGO. In the present study, a red-shift to longer wavelength of the absorption edge of the  $\text{TiO}_2/\text{RGO}$  nanocomposites has also been observed, indicating a narrowing of the band gap of  $\text{TiO}_2$  which should improve charge transfer from the valence band to the conduction band. The red shift has been assigned to the formation of Ti-O-C bonds and the direct interaction between  $\text{TiO}_2$  and RGO [29,30,32,34,41-44], similar to that observed for carbon-doped  $\text{TiO}_2$  nanoparticles [40], and consistent with the results from FTIR spectroscopy and XPS analysis. The maximum red-shift is observed for the G1.00 sample, which also has an absorption tail spanning the visible spectrum because of the formation of nanocomposites.

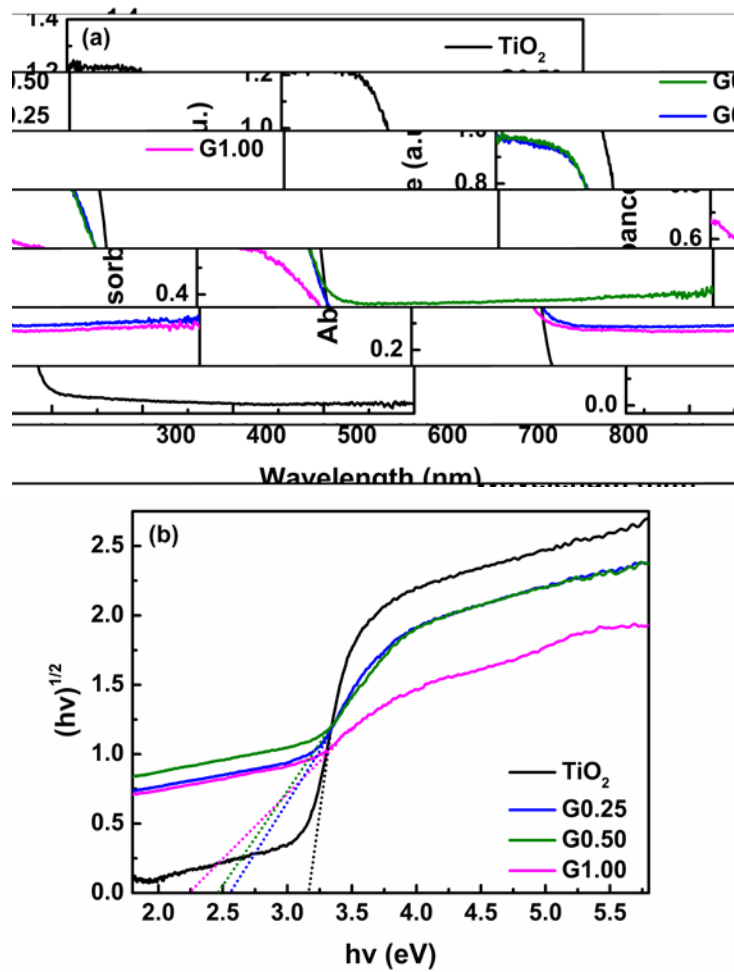


Figure 4. (a) UV-Vis diffuse reflectance spectra of the  $\text{TiO}_2$ , G0.25, G0.50 and G1.00 samples; (b) Tauc plots of the transformed Kubelka-Munk function versus the energy

of light.

The indirect band gap,  $E_g$ , of TiO<sub>2</sub>/RGO nanocomposites near the absorption edge can be obtained from the plot of the modified Kubelka-Munk function, i.e.  $(\alpha h\nu)^{1/2}$ , versus the energy of exciting light ( $h\nu$ ). The absorption coefficient  $\alpha$  is related to  $E_g$  by:

$$\alpha h\nu = A(h\nu - E_g)^{1/2}$$

where  $h\nu$  is the incident photon energy and  $A$  is a constant for indirect transition [44]. Figure 4(b) displays the intercepts of tangents to the  $(\alpha h\nu)^{1/2}$  versus photon energy ( $h\nu$ ) plots, from which the (approximate) band gaps 2.26, 2.46, 2.57 and 3.16 eV for G1.00, G0.50, G0.25 and TiO<sub>2</sub>, respectively, can be obtained. This indicates that the interaction between TiO<sub>2</sub> and RGO increases with the increase in TiO<sub>2</sub> content, and further supports the qualitative observation of a red-shift in the absorption edge of TiO<sub>2</sub>/RGO nanocomposites compared with the pure TiO<sub>2</sub>. The measured energy gap of 3.16 eV for TiO<sub>2</sub> is consistent with the anatase crystalline state. Similar results were also observed for carbon nanotubes/TiO<sub>2</sub> composites due to chemical bonding between the specific carbon sites and TiO<sub>2</sub> [45]. The maximum band-gap narrowing (0.9 eV) is comparable with the value of 0.14 eV observed for carbon-modified TiO<sub>2</sub> [40].

### 3.5 Photoluminescence spectra analysis

For nanostructured materials, photoluminescence spectra are a useful probe of the transfer efficiency of the photoinduced electrons and holes, because the photoluminescence emission comes from recombination of the free charge carriers [46]. Figures S4 and 5 display the excitation and emission photoluminescence spectra of pure TiO<sub>2</sub> and the TiO<sub>2</sub>/RGO composites (G0.25, G0.50 and G1.00). The three TiO<sub>2</sub>/RGO composites exhibit an excitation spectrum centered at *ca.* 348 nm, behavior consistent with that of pure TiO<sub>2</sub> (Figure S4). Upon excitation at 348 nm, the photoluminescence spectra of the TiO<sub>2</sub>/RGO composites are similar to that of pure TiO<sub>2</sub> but with different intensity (Figure 5). It is well known that photoluminescence emission is the result of recombination of the excited holes and electrons either indirectly (via a band-gap state) or directly (band-band), so a higher

photoluminescence intensity may indicate a higher recombination rate of electrons and holes [47,48]. There is a decrease in photoluminescence intensity in proceeding from pure  $\text{TiO}_2$  to the composite samples. This can be attributed to the electrons being excited from the valence band to the conduction band and then transferred to the RGO sheets because of the abundant electron-accepting sites of the latter, hindering electron-hole pair recombination. It can therefore be concluded that the presence of RGO in the nanocomposite samples might be effective in enhancing the separation efficiency of electron-hole pairs. If this is the case, then this photoluminescence quenching may afford a measure of the interfacial charge transfer via the Ti-O-C linkages and consequent electronic interactions between the RGO and  $\text{TiO}_2$  nanoparticles in the  $\text{TiO}_2/\text{RGO}$  nanocomposites [6]. The photoluminescence of these samples was also investigated with an excitation wavelength of 300 nm (Figure S5). Consistent with the aforementioned results, a significant fluorescence quenching was observed for the  $\text{TiO}_2/\text{RGO}$  nanocomposites, providing further evidence for significant electronic communication between  $\text{TiO}_2$  and RGO in the excited state.

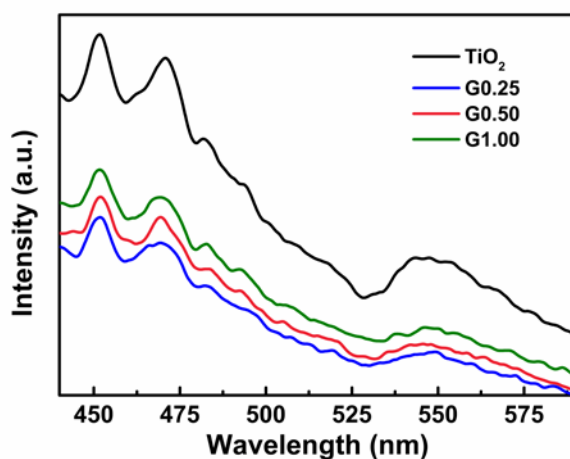


Figure 5. Emission photoluminescence spectra of the  $\text{TiO}_2$ , G0.25, G0.50 and G1.00 samples.

### 3.6 Raman spectra analysis

Raman spectroscopy was used to characterize the  $\text{TiO}_2/\text{RGO}$  nanocomposites as well as the as-prepared RGO. RGO exhibits a prominent D-band at  $1345\text{ cm}^{-1}$ ,

corresponding to the breathing mode of  $\kappa$ -point phonons of  $A_{1g}$  symmetry. There is also an intense tangential mode at  $1580\text{ cm}^{-1}$  (the G-band), which arises from the first-order scattering of  $E_{2g}$  phonons from  $sp^2$  carbon atoms [18], as shown in the inset of Figure 6. For the  $\text{TiO}_2/\text{RGO}$  composites, several characteristic bands were observed at ca. 152, 399, 513, and  $636\text{ cm}^{-1}$ , corresponding to the  $E_{g(1)}$ ,  $B_{1g(1)}$ ,  $A_{1g} + B_{1g(2)}$ , and  $E_{g(2)}$  modes of anatase, respectively [49]. Generally, the full-width at half-maximum (FWHM) of the anatase  $E_g$  mode increases with decreasing crystallite size of anatase [50]; in the present case, the peak intensities of the  $E_g$  mode are significantly increased and the peak widths are broadened on increasing addition of  $\text{TiO}_2$ , consistent with decreasing particle size. Significantly, D- and G-bands for graphitized structures were also observed, confirming the presence of RGO in the  $\text{TiO}_2/\text{RGO}$  composites. The frequencies of these spectral features match those of pristine graphite, consistent with the reduction of GO. The intensity ratio of the D/G bands provides a measure of the disorder/defects in RGO, a smaller intensity ratio of  $I_D/I_G$  corresponding to larger average size and fewer  $sp^3$  defects/disorders of the in-plane graphitic crystallite  $sp^2$  domains. In the present case, the intensity ratio of  $I_D/I_G$  for RGO is 0.94, increasing to 1.06, 1.03 and 1.05 for G0.25, G0.50 and G1.00, respectively. This is consistent with the newly formed graphitic domains being smaller in size and greater in number than those in GO before the hydrothermal treatment, as reported by Shah and coworkers [51]. The plane crystallite size,  $L_a$ , characterizing the RGO domains can be obtained from the equation:

$$L_a(\text{nm}) = \frac{560}{E_\lambda^4} \left( \frac{I_D}{I_G} \right)^{-1}$$

where  $E_\lambda$  is the excitation energy in eV [1]. Thus, RGO has an  $L_a$  of  $\sim 20.2\text{ nm}$ , which upon functionalization reduces to  $\sim 17.9$ ,  $\sim 18.4$  and  $\sim 18.1\text{ nm}$  for G0.25, G0.50 and G1.00, respectively. This change further confirms a decrease in the average size and a higher defect density of the in-plane graphitic crystallite  $sp^2$  domains upon reduction of the exfoliated GO, suggesting a decrease in the oxygen-containing groups on RGO, and consistent with the results of FTIR.

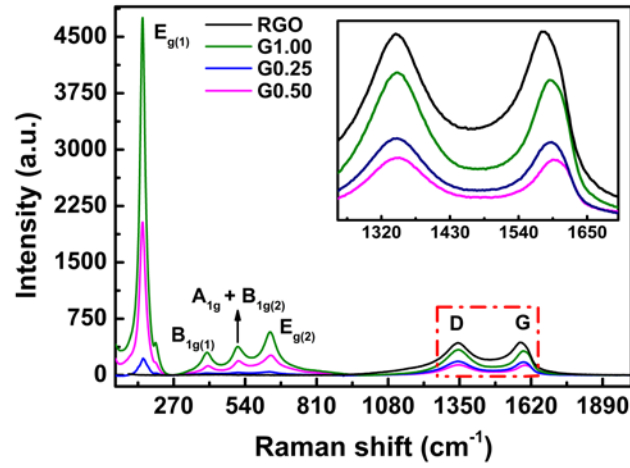


Figure 6. Raman spectra of the RGO, G0.25, G0.50 and G1.00 samples.

### 3.7 XRD studies

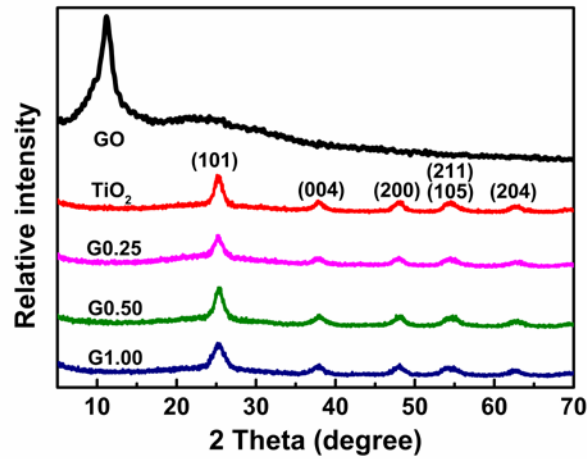


Figure 7. XRD patterns of the GO, TiO<sub>2</sub>, G0.25, G0.50 and G1.00 samples.

Figure 7 displays the XRD patterns of the TiO<sub>2</sub>, G0.25, G0.50 and G1.00 samples, together with the pattern of GO. In the case of GO, the peak at 11.2° indicates a large interlayer distance of 0.78 nm that results from the presence of oxygenated functional groups on the graphene sheets. The reflections in the XRD patterns of TiO<sub>2</sub> and the TiO<sub>2</sub>/RGO composites can be indexed to anatase TiO<sub>2</sub> (JCPDS Card No. 21-1272) [37,52,53]; no rutile phase TiO<sub>2</sub> was detected in this work. The typical diffraction peaks of RGO were not found in the XRD patterns of the nanocomposites, a result



that may derive from the main peak of anatase  $\text{TiO}_2$  masking the main characteristic peak of RGO [24]. The same phenomenon was observed in a previous paper [53], indicating that GO was reduced to RGO during the hydrothermal process, and consistent with the results from XPS. Similar XRD patterns of anatase  $\text{TiO}_2$  nanoparticles from different samples indicate that the incorporation of RGO has little influence on the phase structure of  $\text{TiO}_2$  during the hydrothermal process.

### 3.8 TEM studies

Typical TEM images of the RGO and  $\text{TiO}_2/\text{RGO}$  composites are shown in Figure 8. Samples for TEM analysis were obtained by spreading a drop of a dilute dispersion of the as-prepared products on amorphous carbon-coated copper grids that were then dried in air before transfer to the TEM sample chamber. Figure 8(a) reveals that, as expected, the RGO starting material is composed of single- and multi-layer RGO sheets; a large single-layer RGO flake is visible, as well as a large number of multi-layer RGO sheets. The flakes are almost transparent and entangled with each other, resembling a crumpled silk veil [54]; it is impossible to determine the number of layers. The TEM images of the  $\text{TiO}_2/\text{RGO}$  nanocomposites reveal that the composites are composed of RGO sheets and many small  $\text{TiO}_2$  nanoparticles. Since the monolayer RGO sheets are extremely thin, it was difficult to distinguish between them and the carbon-supported films on the copper grid. Nevertheless, the edges and crumpled silk waves of these RGO nanosheets confirm that the  $\text{TiO}_2$  nanoparticles are indeed deposited on the supports (the almost transparent RGO sheets). During formation of the composites, the  $\text{TiO}_2$  nanoparticles adhere to the functional groups on the GO plane, with simultaneous reduction of GO to RGO by electron donation from ethanol and formation of Ti-O-C bonding between  $\text{TiO}_2$  and RGO [55]. A majority of the  $\text{TiO}_2$  nanoparticles are anchored to the RGO sheets (there is a very small amount of free-standing  $\text{TiO}_2$  particles), which is very important as it ensures efficient electron collection via the RGO sheets during the insertion/extraction processes. Our TEM measurements further reveal that the increase in tetrabutyl titanate content in the hydrothermal process increases the dispersion of the  $\text{TiO}_2$  nanoparticles on the RGO sheets.

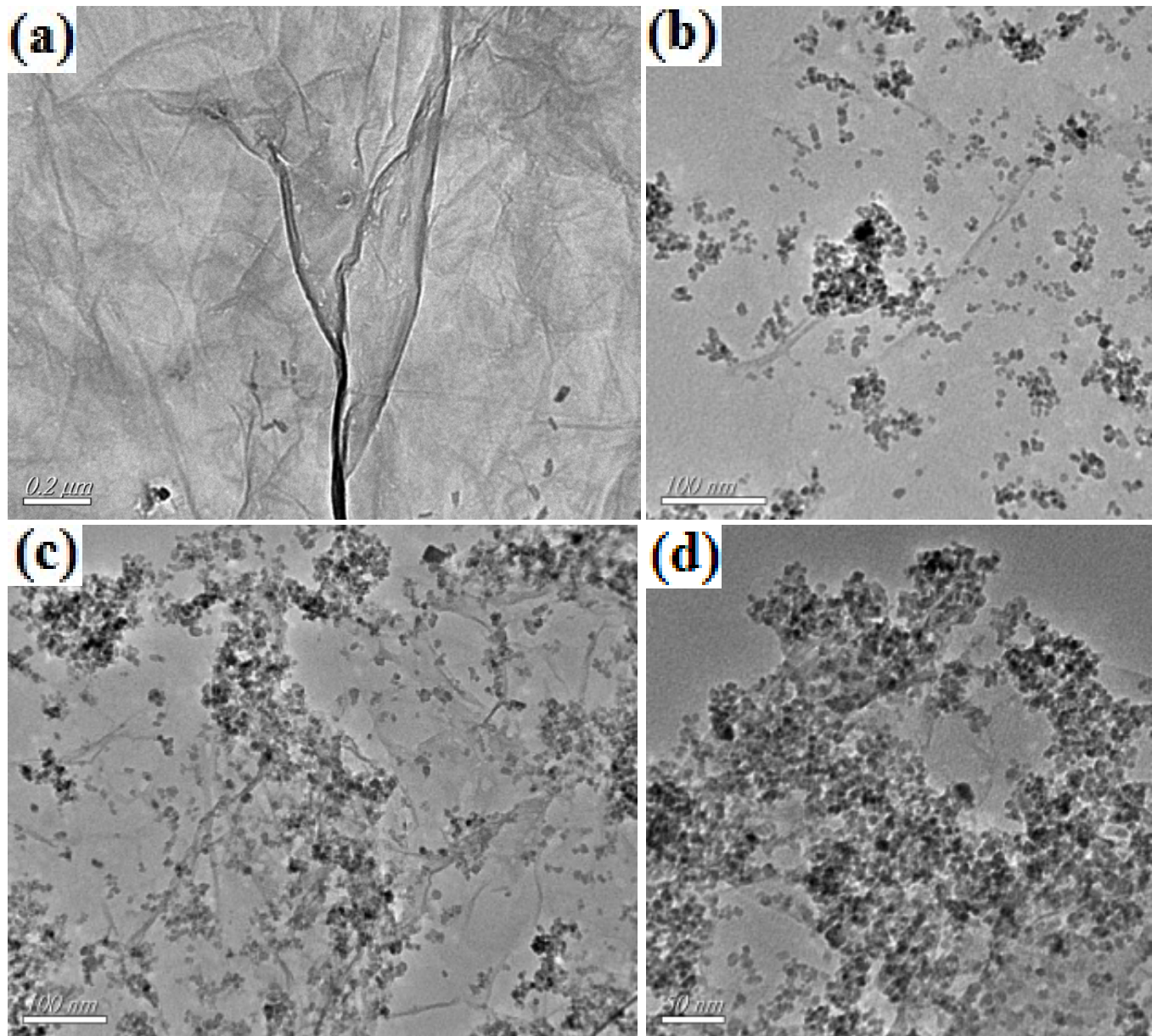


Figure 8. TEM spectra of the RGO (a), G0.25 (b), G0.50 (c) and G1.00 (d) samples.

### 3.9 Nonlinear absorption and optical limiting behaviour

Recent process in the field of photonics and optoelectronics has encouraged the search for materials with excellent NLO properties and ultrafast response, with the goal of protecting sensitive optical systems and human eyes [56]. Various materials such as organic molecules [57], inorganic complexes [58], organometallic compounds [59], nanocomposites, and semiconductors [60] have been intensively studied (both theoretically and experimentally) for potential applications in OL. Among the various possible NLO materials, nanocomposites have proven to be superior candidates

because of their outstanding mechanical, electrical and thermal properties [1,61]. Nanocomposites can exhibit strong nonlinear extinction (including, inter alia, absorption, scattering and/or refraction) for intense laser beams and thus may be viable OL materials [62]. Successful implementation of their exceptional performance is still lacking due to the limited knowledge of the microstructures of nanomaterials, information on which is needed to design and control their physicochemical properties and to clarify their NLO mechanisms. Recently, the NLO properties of  $\text{TiO}_2$  in nanoscale composite materials has received considerable interest (for example, large reverse-saturable absorption resulting from an accumulative effect of two-photon absorption and electric charge interaction was observed in poly(methyl methacrylate)/ $\text{TiO}_2$  nanocomposites [63]), while zinc oxide-decorated RGO sheets have been shown to exhibit excellent nonlinear properties compared to the individual components [6]. However, NLO studies of metal decorated RGO, especially metal oxide-decorated RGO, are still scarce, so the OL properties of the nanocomposites prepared in the present studies have been assessed.

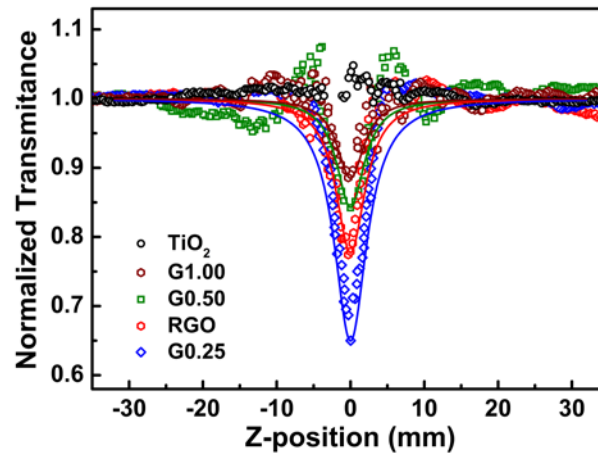


Figure 9. Normalized transmittance curves of pure  $\text{TiO}_2$ , RGO, G0.25, G0.50 and G1.00. Solid curves represent the theoretical fits to the experimental data.

Figure 9 displays representative open-aperture Z-scan curves of pure  $\text{TiO}_2$ , RGO and  $\text{TiO}_2$ -doped RGO nanocomposites as solutions in DMF. For ease of comparison, the concentrations of all samples were adjusted to the same linear transmittance (60%

at 532 nm). As shown in Figure 9, solutions of G0.25, G1.00 and RGO exhibit decreased normalized transmittance with valley-shaped curves as they are brought to the beam waist, consistent with a strong reverse saturable absorption (RSA) response with 4 ns pulse excitation, and OL behavior. For an optical limiter, the depth of the valley in the open-aperture Z-scan curves is correlated with the OL merit [64]. The largest decrease in transmittance value is observed for the G0.25 nanocomposite, suggesting that it should have the best OL performance. Because the solvent DMF exhibits no nonlinear absorption under these experimental conditions, the observed nonlinear absorption must arise from the solutes. The nonlinearity of pure TiO<sub>2</sub> was also investigated under the same experimental conditions, but no obvious signal was observed, suggesting that the TiO<sub>2</sub> has a very small NLO effect that can be neglected. For a quantitative comparison, we have carried out a theoretical fitting of the experimental results to obtain effective nonlinear absorption coefficients  $\beta_{\text{eff}}$  (m/W). With the same linear transmittance, the G0.25 nanocomposite exhibits a better NLO effect, with a nonlinear absorption coefficient  $\beta_{\text{eff}}$  of  $6.0 \times 10^{-10}$  m/W, compared to those of RGO ( $3.0 \times 10^{-10}$  m/W), G1.00 ( $1.3 \times 10^{-10}$  m/W) and G0.50 ( $1.9 \times 10^{-10}$  m/W), indicating a significant influence of the TiO<sub>2</sub> concentration in the nanocomposites on the NLO response. Similar phenomena have been observed with multiwalled carbon nanotube-organically modified silicate nanohybrid gel glass (MWCNT-ORMOSIL) [65]. Comparing the values of  $\beta_{\text{eff}}$  for the G0.25 nanocomposite obtained in the present work with that reported previously for the MWCNT-ORMOSIL system in Ref. [65] reveals that G0.25 has a larger effective nonlinear absorption coefficient, though the differing experimental geometries render such a comparison necessarily cautious.

Interestingly, the curve for G0.50 clearly displays firstly saturable absorption (SA) and then RSA behavior, and then a return to SA behavior as the sample traverses in the  $z$ -direction. Slower excited-state ( $S_1$ ) relaxation and faster higher-lying excited-state ( $S_n$ ) relaxation play a dominant role in switching the NLO absorption performance from SA to RSA and back to SA [66]. The  $S_1$  state is populated by relaxation from the higher excited state and from ground-state absorption. The  $S_n$  state

is populated by ESA corresponding to the  $S_1$ - $S_n$  transition [67]. The lifetime of the  $S_1$  state is longer, and thus G0.50 is readily saturated in the (relatively) low-intensity region. At higher intensities, two-photon absorption (TPA) and the formation of strong light scattering centers may dominate for G0.50, resulting in a RSA effect [68]. Taken together, the observed Z-scan profile is consistent with a combination of ESA (the  $S_1$ - $S_n$  transition), light scattering and a TPA process. Pure RGO displays a clear-cut RSA behavior resulting from a TPA process. The present study reveals that there is an optimum concentration of  $\text{TiO}_2$  in the nanocomposites for NLO behavior, specifically compositions G0.25 and G0.50, which have the strongest RSA behavior and combination of NLO mechanisms, respectively. The nonlinear absorption of nanocomposite G1.00 is the weakest, possibly due to the loading of  $\text{TiO}_2$  being so great that it results in formation of a hydrolyzed titanium alkoxide network structure instead of individual oxide nanoparticles [63]; as a result, the quantum confinement effect responsible for NLO performance decreases significantly. The magnitude of the nonlinear absorption observed for the three nanocomposites is relatively small compared to previously-reported porphyrin-graphene or phthalocyanine-graphene nanocomposites [69-71]. However, the different NLO response obtained with these nanocomposites suggests that  $\text{TiO}_2$  doped-RGO holds promise for various potential applications. For example, the significant NLO performance of the G0.25 nanocomposite may lead to OL applications in eye and sensor protection, while the NLO mechanism changing from SA to RSA and back to SA as a function of light intensity for G0.50 indicates that it can be a potential candidate for ultrafast NLO switching applications.

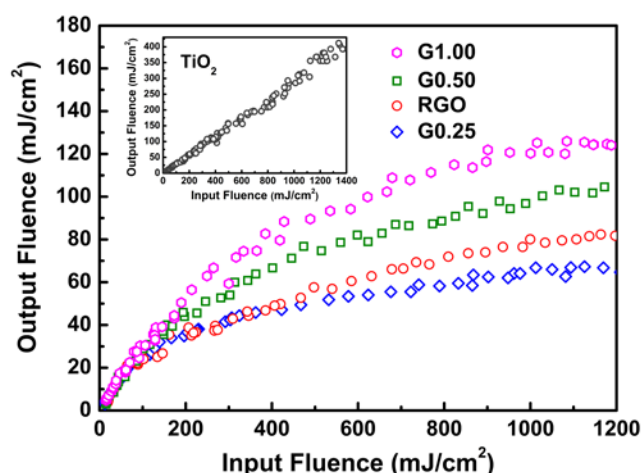


Figure 10. Comparison of the optical limiting performance of RGO, G0.25, G0.50 and G1.00 in DMF at 532 nm. Shown in the inset is the optical limiting response of TiO<sub>2</sub>.

RGO, G0.25, G0.50 and G1.00 exhibit nonlinear absorption, so these materials can in principle function as optical limiters to nanosecond 532 nm laser pulses (Figure 10). For comparison, the OL performance of TiO<sub>2</sub> was investigated, but no obvious limiting response was observed, consistent with the result from open-aperture Z-scan (Figure 9). From Figure 10, it can be clearly seen that, at low input fluences, dispersions of these species (RGO, G0.25, G0.50 and G1.00) obey Beer's Law (linear relationships between input fluence and output fluence). As the incident fluence increases, the output fluence deviates from linearity, consistent with an OL response. Experiments with pure DMF solvent exhibited no detectable OL response under the same conditions, implying that the solvent contribution is negligible; the observed OL effect should thus be assigned solely to the solutes. Compared to the RGO, G0.50 and G1.00 dispersions, the G0.25 nanocomposite displays an enhanced OL response, consistent with the results of open-aperture Z-scan, and presumably due to a combination of the NLO mechanisms of RGO and TiO<sub>2</sub> in the nanocomposite G0.25; RGO exhibits large ESA and TPA in the nanosecond regime, resulting in enhanced nonlinear absorption performance in RGO-containing composites [72]. As shown in Figure 9, G0.25 has the largest dip among the normalized transmittance curves, and thus G0.25 exhibits a much better OL response than RGO, G0.50 or G1.00. The



enhanced OL in G0.25 may have a significant contribution from defect-induced states that are generated during preparation, and which facilitate interband transitions through ESA [1,73]; the increase in the number of defects is confirmed by the increased  $I_D/I_G$  ratio after heat-treatment (Figure 6). Significant fluorescence quenching of  $TiO_2$  following doping with RGO is observed, due to the electron/energy transfer from  $TiO_2$  to RGO. This doping may inhibit the recombination of electron-hole pairs, and thus accelerate the transfer of photogenerated electrons. During the synthesis of the G0.25 nanocomposite, some oxygen-containing groups on the surface of GO are partially removed by hydrothermal treatment, and the extended conjugation is thereby further increased. The larger the extent of conjugation, the higher the chance for electron/energy transfer, leading to enhanced OL response [6].

#### 4. Conclusion

In summary, we have reported the hydrothermal synthesis and structural, spectroscopic and NLO properties of  $TiO_2$ /RGO nanocomposites. The formation of the nanocomposites was confirmed by FTIR, XPS, optical absorption, XRD and Raman spectroscopies, and TEM micrographs. Compared to  $TiO_2$ , significant fluorescence quenching was observed for all nanocomposites, suggesting energy transfer. The NLO absorption properties and optical limiting responses of the nanocomposites were investigated using the Z-scan technique at 532 nm with 4 ns laser pulses. The present study suggests that there is an optimum concentration of  $TiO_2$  in the nanocomposites to maximize NLO behavior; the G0.25 nanocomposite exhibits enhanced NLO performance in the nanosecond regime, ascribed to a combination of NLO mechanisms, while intensity-dependent behavior in G0.50 is observed because of the combined contribution of ESA and TPA. Overall, these results suggest that  $TiO_2$ /RGO nanocomposites with an appropriate content of  $TiO_2$  afford good candidates for optical limiting and NLO switching, and potentially attractive possibilities for ultrafast nonlinear optical materials for photonic and optoelectronic devices.

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### Figure captions

Figure 1. Illustration of the fabrication process of the TiO<sub>2</sub>-decorated RGO nanocomposites.

Figure 2. FTIR spectra of TiO<sub>2</sub>, GO, RGO, G0.25, G0.50 and G1.00. Shown in the inset is the deconvolution spectra of G0.25, G0.50 and G1.00.

Figure 3. (a) XPS survey spectra of GO, RGO, G0.25, G0.50 and G1.00; (b) Ti2p core level XPS spectra of TiO<sub>2</sub> and G0.25; High-resolution XPS spectra of C1s for (c) GO and (d) the G0.25 nanocomposite.

Figure 4. (a) UV-Vis diffuse reflectance spectra of the TiO<sub>2</sub>, G0.25, G0.50 and G1.00 samples; (b) Tauc plots of the transformed Kubelka-Munk function versus the energy of light.

Figure 5. Emission photoluminescence spectra of the TiO<sub>2</sub>, G0.25, G0.50 and G1.00 samples.

Figure 6. Raman spectra of the RGO, G0.25, G0.50 and G1.00 samples.

Figure 7. XRD patterns of the GO, TiO<sub>2</sub>, G0.25, G0.50 and G1.00 samples.

Figure 8. TEM spectra of the RGO (a), G0.25 (b), G0.50 (c) and G1.00 (d) samples.

Figure 9. Normalized transmittance curves of pure TiO<sub>2</sub>, RGO, G0.25, G0.50 and G1.00. Solid curves represent the theoretical fits to the experimental data.

Figure 10. Comparison of the optical limiting performance of RGO, G0.25, G0.50 and G1.00 in DMF at 532 nm. Shown in the inset is the optical limiting response of TiO<sub>2</sub>.